# Prediction and Analysis of the Force and Shape Parameters in Variable Gauge Rolling 

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#### Abstract

Variable gauge rolling is a new process to obtain a plate for which the thickness changes continuously by continuously and dynamically adjusting the roll gap upward and downward in the rolling process. This technology is an effective method for producing lightweight, low-cost, and economical plates. However, variable gauge rolling is an unsteady process, and the changes in the force and deformation parameters are complex. In this research, based on the minimum energy theory of the variational principle and considering the characteristics of the roll movement and workpiece deformation comprehensively, the internal plastic deformation, friction, shear and tension powers, and the minimum result of the total power functional in upward and downward rolling are obtained with the first integral and then with a variation of adopting the specific plastic power and strain rate vector inner product. The analytical results of the deformation and force parameters are also established using the variational method. Then the precision of this model is certified using the measured values in a medium plate hot rolling plant and the experimental data for Tailor Rolled Blank rolling. Good agreement is found. Additionally, the variation rule of bite angle, neutral angle, and location neutral points are shown, and the change mechanism of the friction parameter on the stress state effect coefficient is given in variable gauge rolling. This research proposes a new mathematical model for rolling process control that provides a scientific basis and technical support for obtaining an accurate section shape in variable gauge rolling production.


Keywords: Variable gauge rolling, Analytical solution, Roll separating force, Energy method, Neutral point

## 1 Introduction

Variable gauge rolling is applied to produce plates that have the benefit of being lightweight due to the optimization of the plate thickness according to the load distribution. If traditional equal-thickness steel plates are replaced by variable thickness plates with comparable properties, more than $30 \%$ of the steel in the construction of bridges (such as the steel in the main beam flange, and web) can be saved [1], and more than $40 \%$ of the weight in the manufacture of automobile parts (such as reinforcement beams, and instrument panel bottom crossbeams)

[^0]can be reduced [2]. Variable thickness plates obtained by rolling have the advantages of decreased frame weights, reduction of welds, high production efficiency, and yield ratios [3]. Variable thickness plates have been preliminarily applied in automobile manufacturing, construction, shipbuilding bridge industries and so on [4]. The demand for variable thickness plates is growing [5].
There are two different transition zones from the thick area to the thin area and from the thin area to the thick area. The shape and the size of these transition zones not only determine the bearing capacity during service, but also influence the subsequent processing and die design [6]. Nevertheless, the high prediction accuracy force parameters (such as the roll torque, and roll separating force) and deformation parameters (such as the location of the neutral point, and forward slip) for different
processing parameters are the keys to obtaining excellent shapes for the transition zones in variable thickness plates [7, 8].
Variable gauge rolling has been widely used in Mizushima Automatic Plan View Pattern Control System rolling (MAS rolling) [9], Tailor Rolled Blank (TRB) [10], Longitude Profile (LP) hot rolling [11], and Flying Gauge Change (FGC) [12]. To compensate for the irregular shape of the workpiece in medium plate rolling, the MAS rolling method is adopted to improve the shape in the middle pass. Then the medium plate not only reduces the loss of the head, tail, and edges, but it also enhances the metal yield and the yield ratio [13], as shown in Figure 1(a). TRB is used to replace Tailor Welded Blank (TWB), which is obtained by welding two or more plates of different thicknesses together [14]. TRBs are obtained by directly cutting a periodic variable thickness plate produced through flexible rolling technology, as shown in Figure 1(b). TRB has apparent features such as no weld joint, good surface quality, reliable subsequent processing, and high production efficiency, suitability for large-scale production, low production cost, and energy savings [15]. TRB adopts a new processing technology to reduce the weight of structural plates, and meets the performance requirements of the different parts of a vehicle.


Figure 1 Schematic diagram of typical variable gauge rolling

Therefore, it reduces the bodyweight of vehicles and can be used to research and develop environmentally friendly and energy-saving automobiles.
Ding et al. [16] proposed a method for accurately analyzing and optimizing the shape of the transition zone in TRB based on the isogeometric analysis method. According to the rolling deformation features, Zhi et al. [17] designed four kinds of TRB transition curve functional equations, namely, double arc, straight line, concave arc, and power function equations, and analyzed the roll center trajectory during variable gauge rolling. Wang et al. [18] proposed the design method of transition zone curve, and deduced the velocity models of workpiece's horizontal motion and upward and downward roll motion, but the specific procedure to calculate the neutral angle contained in these velocity models was not given. Wang et al. [12] adopted variable gauge rolling technology in FGC and studied the setting principle of the roll gap and roll circumferential speed. Yu et al. [19] derived the forward slip equation in the variable gauge rolling process according to the equal flow per second, and analyzed the influence of different kinds of transition zone shapes on the forward slip, and compared them with the results of a finite element simulation. Because the deformation characteristics outside the transition zone in variable gauge rolling are similar to the characteristics in traditional rolling, Zhang et al. [20] calculated the rolling speeds in the thick area and the thin area with the traditional rolling theory, and then took the mean value of the two rolling speeds as the speed of the transition zone.
Liu [1] systematically researched the theory and application of variable gauge rolling. The bite condition and the length of the contact arc were analyzed first, and then the formulas of the forward slip and neutral angle were deduced based on the assumption that the unit pressure was uniformly distributed along the contact arc. The changes in the roll separating force were analyzed using the finite element method in variable gauge rolling. Zhang [21] established force equilibrium differential equations in TRB rolling with the slab method, and the unit pressure distribution in the deformation zone was obtained by referring to the idea of solving the Karman differential equation. Then the calculation formulas of the roll separating force, neutral angle, and forward and backward slips were determined. Yu et al. [22] analyzed the fluctuation of roll separating force in variable gauge rolling using finite element software, and researched the influences of the reduction, friction, and radius on the roll separating force.

The exit position of the workpiece is not on the centerline of the two cylindrical work rolls in the variable gauge rolling, so the thickness of the exit zone is not equal to the loaded roll gap. The variable gauge rolling process cannot be identical to the combination of the infinite equal thickness traditional rolling process, and the velocity fields of the two processes are different even if the instantaneous inlet and outlet thicknesses are the same [23]. Therefore, in this research, a velocity field satisfying kinematically admissible conditions is established based on the mass conservation law in the deformation zone with consideration of the influence of the upward and downward motions of the rolls. The high accuracy mechanical and deformation parameters in upward and downward rolling are obtained using suitable mathematical solution methods.

## 2 Velocity Field

The Cartesian coordinate system is located at the central node of the centerline of the two work rolls. The length, thickness, and width directions of the workpiece are

(b) Finished product

Figure 2 Schematic diagram of the deformation zone and finished product in upward rolling


Figure 3 Schematic diagram of the deformation zone and finished product in downward rolling
expressed by the $x, y$, and $z$ axes separately, as shown in Figures 2 and 3. For the work rolls, the original radius is represented by $R_{0}$, and the flattened radius is represented by $R$. The workpiece thickness of the thick zone decreases from $2 h_{0}$ to $2 h_{\text {thick }}$, and the thickness of the thin zone is reduced from $2 h_{0}$ to $2 h_{\text {thin }}$. The exit thickness of the deformation zone at some point is $2 h_{1}$, so the unilateral reduction is $\Delta h=h_{0}-h_{1} . v_{R}$ is the roll circumferential speed. $V_{y}$ is the speed of the upward or downward motion of the rolls. $\theta$ is the bite angle, and $\alpha$ is the contact angle. $\alpha_{0}$ is the wedge angle of the variable thickness section, and $x_{0}$ is the distance from the exit to the centerline of the roll.
According to the directions of the Cartesian coordinate system, $V_{y}$ is a positive value in the upward rolling, and it is a negative value in the downward rolling. To simplify the calculation, this model ignores the influence of spreading, and the width of the workpiece is $2 w$. A quarter deformation zone is used to build the model to simplify the calculation due to the symmetry. The expressions for half of the thickness $h_{x}\left(h_{\alpha}\right)$ of the deformation zone and the first derivative $h_{x}^{\prime}$ are described as Eqs. (1)-(4).

$$
\begin{align*}
& h_{x}=h_{0}+\sqrt{R^{2}-l^{2}}-\sqrt{R^{2}-x^{2}}  \tag{1}\\
& h_{\alpha}=h_{0}+R \cos \theta-R \sin \alpha  \tag{2}\\
& x=R \sin \alpha, \mathrm{~d} x=R \cos \alpha \mathrm{~d} \alpha \tag{3}
\end{align*}
$$

$$
\begin{equation*}
h_{x}^{\prime}=\frac{x}{\sqrt{R^{2}-x^{2}}}=\tan \alpha \tag{4}
\end{equation*}
$$

Considering the influence of the upward or downward motions of the rolls during variable gauge rolling, the constant volume condition of the metal flow satisfies

$$
\begin{equation*}
h_{0} v_{0}=h_{x} v_{x}+V_{y} R(\theta-\alpha) \tag{5}
\end{equation*}
$$

The velocity components during variable gauge rolling are

$$
\left\{\begin{align*}
v_{x} & =\frac{h_{0} v_{0}-V_{y} R \theta+V_{y} R \alpha}{h_{x}}  \tag{6}\\
v_{y} & =\left(\frac{v_{x}}{h_{x}} h_{x}^{\prime}-\frac{V_{y} R}{x h_{x}} h_{x}^{\prime}\right) y \\
v_{z} & =0
\end{align*}\right.
$$

The strain rate field components during variable gauge
meets the constant volume condition. Eqs. (6) and (7) fulfill the kinematically admissible conditions.

## 3 Analytical Model of Total Power Functional in Upward Rolling

Because the deformation zones in variable gauge and traditional rolling are different, the expressions for the length of the contact arc $l_{\text {total }}$ and the thickness of the exit zone $h_{1}$ in the deformation zone during the upward rolling according to Figure 2 are

$$
\begin{equation*}
l_{\mathrm{total}}=l-x_{0}=\sqrt{R^{2}+\left(R \cos \alpha_{0}-\Delta h\right)^{2}}-x_{0} \tag{8}
\end{equation*}
$$

$$
\begin{equation*}
h_{1}=h_{\mathrm{hthin}}+V_{y} t \tag{9}
\end{equation*}
$$

where $t$ is the time of the variable gauge rolling.

### 3.1 Internal Plastic Deformation Power

To solve the calculation for the internal plastic deformation power difficulty that is caused by the nonlinear Mises yield criterion, the internal plastic deformation power can be received using the MY liner yield criterion in upward rolling. The MY liner yield criterion is the arithmetic mean of the twin shear stress and Tresca yield criterion loci [24]. Noting that $\dot{\varepsilon}_{\max }=\dot{\varepsilon}_{x}, \dot{\varepsilon}_{\min }=\dot{\varepsilon}_{y}$ in Eq. (7) and the internal plastic deformation power $\dot{W}_{u i}$ is

$$
\begin{align*}
\dot{W}_{\mathrm{ui}} & =\int_{V} D\left(\dot{\varepsilon}_{i j}\right) \mathrm{d} V=\frac{16}{7} \sigma_{\mathrm{s}} \int_{x_{0}}^{l} \int_{0}^{w} \int_{0}^{h_{x}}\left(\dot{\varepsilon}_{\max }-\dot{\varepsilon}_{\min }\right) \mathrm{d} x \mathrm{~d} y \mathrm{~d} z  \tag{10}\\
& =-\frac{32}{7} \sigma_{\mathrm{s}} w \int_{x_{0}}^{l}\left[\left(-h_{0} v_{0}+V_{y} R \theta\right) \frac{h_{x}^{\prime}}{h_{x}}-V_{y} R \frac{\alpha}{h_{x}} h_{x}^{\prime}+V_{y} R \frac{h_{x}^{\prime}}{x}\right] \mathrm{d} x .
\end{align*}
$$

rolling are

$$
\left\{\begin{array}{l}
\dot{\varepsilon}_{x}=-\frac{v_{x}}{h_{x}} h_{x}^{\prime}+\frac{V_{y} R}{x h_{x}} h_{x}^{\prime}  \tag{7}\\
\dot{\varepsilon}_{y}=\frac{v_{x}}{h_{x}} h_{x}^{\prime}-\frac{V_{y} R}{x h_{x}} h_{x}^{\prime} \\
\dot{\varepsilon}_{z}=0
\end{array}\right.
$$

According to Eqs. (6) and (7), at the inlet of the zone: $\left.v_{x}\right|_{x=l}=v_{0},\left.v_{y}\right|_{y=0}=0$, and $\left.v_{z}\right|_{z=0}=0$; and at the exit of the zone: $\left.v_{y}\right|_{x=x_{0}} \neq 0$. This is caused by the upward or downward motions of the rolls. The strain rate field satisfies $\dot{\varepsilon}_{x}+\dot{\varepsilon}_{y}+\dot{\varepsilon}_{z}=0$, so the velocity field satisfies the velocity boundary condition, and the strain rate field

To simplify the calculation and writing, it is assumed that $M=h_{0} / R+\cos \theta$. The metal flow volume per second is $U=v_{0} h_{0} w=v_{n} h_{n} w$ $=v_{R} \cos \alpha_{N} w\left(h_{0}+R \cos \theta-R \cos \alpha_{n}\right)$. Therefore, the mean value of $\cos \alpha$ obtained with the integral mean value theorem is

$$
\begin{equation*}
\overline{\cos \alpha}=\frac{\int_{\alpha_{0}}^{\theta} \cos \alpha \mathrm{d} \alpha}{\theta-\alpha_{0}}=\frac{\sin \theta-\sin \alpha_{0}}{\theta-\alpha_{0}} \tag{11}
\end{equation*}
$$

The internal plastic deformation power $\dot{W}_{\mathrm{ui}}$ is finally obtained as

$$
\begin{align*}
\dot{W}_{\mathrm{ui}}= & \frac{32}{7} \sigma_{\mathrm{s}} U \ln \frac{h_{0}}{h_{1}} \\
& +\frac{32}{7} \sigma_{\mathrm{s}} w V_{y} R\left\{\theta \ln \frac{h_{1}(M-\cos \theta)}{h_{0}}\right.  \tag{12}\\
& -\alpha_{0} \ln \left(M-\cos \alpha_{0}\right)-\left(\theta-\alpha_{0}\right) \\
& {\left.\left[1+\ln \left(M-\frac{\sin \theta-\sin \alpha_{0}}{\theta-\alpha_{0}}\right)\right]\right\} . }
\end{align*}
$$

### 3.2 Shear Power

According to the velocity field in Eq. (6), the velocity discontinuity at the inlet of the zone $(x=l)$ is

$$
\begin{align*}
& \dot{W}_{\mathrm{us}}=\dot{W}_{\mathrm{us} 1}+\dot{W}_{\mathrm{us} 2}=2 k U\left(\tan \theta+\tan \alpha_{0}\right) \\
& -\frac{2 k w h_{0} V_{y} R \tan \theta}{l}-2 k w V_{y} R \tan \alpha_{0}\left(\frac{h_{1}}{x_{0}}+\theta-\alpha_{0}\right) . \tag{17}
\end{align*}
$$

### 3.3 Friction Power

The friction only exists on the contact surface between the roll and the workpiece, as indicated in Figure 2. According to Eq. (6), the tangential velocity discontinuity $\Delta v_{f}$ on the contact surface along the $x, y$, and $z$ axes is $\Delta v_{x}, \Delta v_{y}$, and $\Delta v_{z}$, respectively.

$$
\left\{\begin{array}{l}
\Delta v_{x}=v_{R} \cos \alpha-\left(\frac{h_{0} v_{0}-V_{y} R \theta}{h_{x}}+\frac{V_{y} R \alpha}{h_{x}}\right)  \tag{18}\\
\left.\Delta v_{y}\right|_{y=h_{x}}=v_{R} \sin \alpha-V_{y}-\left(\frac{h_{0} v_{0}-V_{y} R \theta}{h_{x}}+\frac{V_{y} R \alpha}{h_{x}}-\frac{V_{y} R}{x}\right) \tan \alpha \\
\Delta v_{z}=0
\end{array}\right.
$$

$$
\left\{\begin{array}{l}
\left.v_{y}\right|_{x=l}=\left(\frac{v_{0}}{h_{0}}-\frac{V_{y} R}{l h_{0}}\right) \tan \theta y \\
\left.v_{z}\right|_{x=l}=0
\end{array}\right.
$$

Therefore, the shear power $\dot{W}_{\text {us } 1}$ at the inlet of the zone is

$$
\begin{align*}
\dot{W}_{\mathrm{us} 1} & =4 k \int_{0}^{h_{0}} \int_{0}^{w}\left|\Delta v_{\mathrm{t}}\right| \mathrm{d} y \mathrm{~d} z \\
& =4 k \int_{0}^{h_{0}} \int_{0}^{w}\left(\frac{v_{0}}{h_{0}}-\frac{V_{y} R}{l h_{0}}\right) \tan \theta y \mathrm{~d} y \mathrm{~d} z  \tag{14}\\
& =2 k w h_{0} \tan \theta\left(v_{0}-\frac{V_{y} R}{l}\right) \\
& =2 k U \tan \theta-\frac{2 k w h_{0} V_{y} R \tan \theta}{l}
\end{align*}
$$

At the exit of the zone $\left(x=x_{0}\right)$

$$
\left\{\begin{array}{l}
\left.v_{y}\right|_{x=x_{0}}=\left(\frac{h_{0} v_{0}-V_{y} R \theta+V_{y} R \alpha_{0}}{h_{1}^{2}}-\frac{V_{y} R}{x_{0} h_{1}}\right) \tan \alpha_{0} y,  \tag{15}\\
\left.v_{z}\right|_{x=x_{0}}=0
\end{array}\right.
$$

Hence, the shear power $\dot{W}_{\text {us2 }}$ at the exit of the zone is

$$
\begin{align*}
\dot{W}_{\mathrm{us} 2} & =4 k \int_{0}^{h_{1}} \int_{0}^{w}\left|\left(\frac{h_{0} v_{0}-V_{y} R \theta+V_{y} R \alpha_{0}}{h_{1}^{2}}-\frac{V_{y} R}{x_{0} h_{1}}\right) \tan \alpha_{0} y\right| \mathrm{d} y \mathrm{~d} z  \tag{16}\\
& =2 k U \tan \alpha_{0}-2 k w V_{y} R \tan \alpha_{0}\left(\frac{h_{1}}{x_{0}}+\theta-\alpha_{0}\right)
\end{align*}
$$

Then the total shear power is

$$
\begin{align*}
& \dot{W}_{\mathrm{uf}}=4 m k w\left(\int_{x_{0}}^{l} \Delta v_{x} \cos \alpha \sec \alpha \mathrm{~d} x+\int_{x_{0}}^{l} \Delta v_{y} \sin \alpha \sec \alpha \mathrm{~d} x\right) \\
& =4 m k w\left\{\int_{x_{0}}^{l}\left[v_{R} \cos \alpha-\left(\frac{h_{0} v_{0}-V_{y} R \theta}{h_{x}}+\frac{V_{y} R \alpha}{h_{x}}\right)\right] \mathrm{d} x\right.  \tag{21}\\
& \left.+\int_{x_{0}}^{l}\left[v_{R} \sin \alpha-V_{y}-\left(\frac{h_{0} v_{0}-V_{y} R \theta+V_{y} R \alpha}{h_{x}}-\frac{V_{y} R}{x}\right) \tan \alpha\right] \tan \alpha \mathrm{d} x\right\} \\
& =4 m k w\left(\int_{x_{0}}^{l} \mathrm{~B}_{\mathrm{I} 1} \mathrm{~d} x+\int_{x_{0}}^{l} \mathrm{~B}_{\mathrm{I} 2} \mathrm{~d} x\right)=4 m k w\left(\mathrm{I}_{1}+\mathrm{I}_{2}\right),
\end{align*}
$$

where $I_{1}$ is

$$
\begin{align*}
\mathrm{I}_{1}= & \int_{x_{0}}^{l} \mathrm{~B}_{\mathrm{I} 1} \mathrm{~d} x=-\int_{x_{0}}^{x_{n}} \mathrm{~B}_{\mathrm{I} 1} \mathrm{~d} x+\int_{x_{n}}^{l} \mathrm{~B}_{\mathrm{I} 1} \mathrm{~d} x \\
= & v_{R} R\left(\frac{\alpha_{0}}{2}+\frac{\theta}{2}-\alpha_{n}+\frac{\sin 2 \theta}{4}-\frac{\sin 2 \alpha_{n}}{2}+\frac{\sin 2 \alpha_{0}}{4}\right)+\left(h_{0} v_{0}-V_{y} R \theta+\right.  \tag{22}\\
& \left.V_{y} R \alpha_{m}\right)\left\langle\alpha_{0}+\theta-2 \alpha_{n}+\frac{2 M}{\sqrt{M^{2}-1}}\left\{2 \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{n}}{2}\right)\right]-\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{0}}{2}\right)\right]-\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\theta}{2}\right)\right]\right\}\right\rangle,
\end{align*}
$$

where the mean value of the contact angle $\alpha_{m}$ during upward rolling is obtained with the integral mean value theorem

$$
\begin{equation*}
\alpha_{m}=\frac{\int_{x_{0}}^{l} \alpha \mathrm{~d} x}{l-x_{0}}=\frac{\theta \sin \theta+\cos \theta-\alpha_{0} \sin \alpha_{0}-\cos \alpha_{0}}{\sin \theta-\sin \alpha_{0}} \tag{23}
\end{equation*}
$$

## Similarly,

$$
\begin{align*}
& \mathrm{I}_{2}=\int_{x_{0}}^{l} \mathrm{~B}_{\mathrm{I} 2} \mathrm{~d} x=-\int_{x_{0}}^{x_{n}} \mathrm{~B}_{\mathrm{I} 2} \mathrm{~d} x+\int_{x_{n}}^{l} \mathrm{~B}_{\mathrm{I} 2} \mathrm{~d} x \\
& =v_{R} R\left(\frac{\alpha_{0}}{2}+\frac{\theta}{2}-\alpha_{n}+\frac{\sin 2 \alpha_{n}}{2}-\frac{\sin 2 \theta}{4}-\frac{\sin 2 \alpha_{0}}{4}\right)+V_{y} R \ln \frac{\cos ^{2} \alpha_{n}}{\cos \alpha_{0} \cos \theta} \\
& +V_{y} R\left(\cos \alpha_{0}+\cos \theta-2 \cos \alpha_{n}\right)+\left(h_{0} v_{0}-V_{y} R \theta+V_{y} R \alpha_{m}\right) \\
& \left\langle 2 \alpha_{n}-\alpha_{0}-\theta+\frac{1}{M} \ln \frac{\tan ^{2}\left(\pi / 4+\alpha_{n} / 2\right)}{\tan \left(\pi / 4+\alpha_{0} / 2\right) \tan (\pi / 4+\theta / 2)}+\frac{2 \sqrt{M^{2}-1}}{M}\right.  \tag{24}\\
& \left\{\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{0}}{2}\right)\right]-2 \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{n}}{2}\right)\right]\right. \\
& \left.\left.+\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\theta}{2}\right)\right]\right\}\right\rangle .
\end{align*}
$$

Substituting Eqs. (22) and (24) into Eq. (21) yields the friction power

$$
\begin{aligned}
& \dot{W}_{\mathrm{uf}}=4 m k w\left\langle v_{R} R\left(\alpha_{0}+\theta-2 \alpha_{n}\right)+V_{y} R \ln \frac{\cos ^{2} \alpha_{n}}{\cos \alpha_{0} \cos \theta}\right. \\
& +V_{y} R\left(\cos \alpha_{0}+\cos \theta-2 \cos \alpha_{n}\right)+\frac{2\left(U / w-V_{y} R \theta+V_{y} R \alpha_{m}\right)}{M \sqrt{M^{2}-1}} \\
& \left\{2 \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{n}}{2}\right)\right]-\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\theta}{2}\right)\right]\right. \\
& -\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{0}}{2}\right)\right] \\
& \left.\left.+\frac{\sqrt{M^{2}-1}}{2} \ln \frac{\tan ^{2}\left(\pi / 4+\alpha_{n} / 2\right)}{\tan \left(\pi / 4+\alpha_{0} / 2\right) \tan (\pi / 4+\theta / 2)}\right\}\right\rangle
\end{aligned}
$$

### 3.4 Tension Power

If there are front tension $\sigma_{\mathrm{f}}$ and back tension $\sigma_{\mathrm{b}}$ effects on the workpiece, the tension power of the deformation zone is

$$
\begin{equation*}
\dot{W}_{\mathrm{uT}}=4\left(\sigma_{\mathrm{b}} w h_{0} v_{0}-\sigma_{\mathrm{f}} w h_{1} v_{1}\right)=4 U\left(\sigma_{\mathrm{b}}-\sigma_{\mathrm{f}}\right) \tag{26}
\end{equation*}
$$

### 3.5 Total Power Functional

According to Eqs. (12), (17), (25) and (26), the total power functional model in the upward rolling based on the first

$$
\begin{equation*}
\frac{\mathrm{d} J_{\mathrm{u}}^{*}}{\mathrm{~d} \alpha_{n}}=\frac{\mathrm{d} \dot{W}_{\mathrm{ui}}}{\mathrm{~d} \alpha_{n}}+\frac{\mathrm{d} \dot{W}_{\mathrm{us}}}{\mathrm{~d} \alpha_{n}}+\frac{\mathrm{d} \dot{W}_{\mathrm{uf}}}{\mathrm{~d} \alpha_{n}}+\frac{\mathrm{d} \dot{W}_{\mathrm{uT}}}{\mathrm{~d} \alpha_{n}}=0 \tag{28}
\end{equation*}
$$

where

$$
\begin{equation*}
\frac{\mathrm{d} \dot{W}_{\mathrm{ui}}}{\mathrm{~d} \alpha_{n}}=\frac{32}{7} \sigma_{\mathrm{s}} N \ln \frac{h_{0}}{h_{1}} \tag{29}
\end{equation*}
$$

$$
\begin{equation*}
\frac{\mathrm{d} \dot{W}_{\mathrm{us}}}{\mathrm{~d} \alpha_{n}}=2 k N\left(\tan \theta+\tan \alpha_{0}\right) \tag{30}
\end{equation*}
$$

$$
\begin{align*}
& \frac{\mathrm{d} \dot{W}_{\mathrm{uf}}}{\mathrm{~d} \alpha_{n}}=8 m k w\left\langle-v_{R} R-V_{y} R \tan \alpha_{n}+\frac{U / w-V_{y} R \theta+V_{y} R \alpha_{m}}{\cos \alpha_{n}\left(M-\cos \alpha_{n}\right)}\right. \\
& +V_{y} R \sin \alpha_{n}+\frac{N}{w M \sqrt{M^{2}-1}}\left\{2 \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{n}}{2}\right)\right]\right.  \tag{31}\\
& -\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{0}}{2}\right)\right]-\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\theta}{2}\right)\right] \\
& \left.\left.+\frac{\sqrt{M^{2}-1}}{2} \ln \frac{\tan ^{2}\left(\pi / 4+\alpha_{n} / 2\right)}{\tan \left(\pi / 4+\alpha_{0} / 2\right) \tan (\pi / 4+\theta / 2)}\right\}\right\rangle .
\end{align*}
$$

variation principle of a rigid-plastic material is defined as Eq. (27) [27].

$$
\begin{equation*}
J_{\mathrm{u}}^{*}=\dot{W}_{\mathrm{ui}}+\dot{W}_{\mathrm{us}}+\dot{W}_{\mathrm{uf}}+\dot{W}_{\mathrm{uT}} \tag{27}
\end{equation*}
$$

By solving the differential of the total power functional $J_{\mathrm{u}}^{*}$ in Eq. (27) with respect to the arbitrary variable $\alpha_{n}$ and making it equal to zero [28], the following equation can be obtained

$$
\begin{equation*}
\frac{\mathrm{d} \dot{W}_{\mathrm{uT}}}{\mathrm{~d} \alpha_{n}}=4 N\left(\sigma_{\mathrm{b}}-\sigma_{\mathrm{f}}\right) \tag{32}
\end{equation*}
$$

where $N=\frac{\mathrm{d} U}{\mathrm{~d} \alpha_{n}}=v_{R} w R \sin 2 \alpha_{n}-v_{R} w\left(h_{0}+R \cos \theta\right) \sin \alpha_{n}$.

## 4 Analytical Model of Total Power Functional in Downward Rolling

According to Figure 3, in the downward rolling process, the expressions for the length of the contact arc $l_{\text {total }}$ and the thickness of the exit zone $h_{1}$ in the deformation zone are defined as Eqs. (33) and (34).

$$
\begin{align*}
& l_{\text {total }}=l+x_{0}=\sqrt{R^{2}-\left(R \cos \alpha_{0}-\Delta h\right)^{2}}+x_{0}  \tag{33}\\
& h_{1}=h_{\text {hthick }}+V_{y} t \tag{34}
\end{align*}
$$

The contact arc in the downward rolling is different from that in the upward rolling, and the deformation zone's length increases in the downward rolling. However, the method used to calculate those powers in the upward rolling can be used in the downward rolling. The internal plastic deformation power can be similarly determined

The shear powers at the inlet and exit of the zones are calculated, and then the total shear power is

$$
\begin{align*}
\dot{W}_{\mathrm{ds}}= & 2 k U\left(\tan \theta+\tan \alpha_{0}\right)-\frac{2 k w h_{0} V_{y} R \tan \theta}{l} \\
& -2 k w V_{y} R \tan \alpha_{0}\left(-\frac{h_{1}}{x_{0}}+\theta+\alpha_{0}\right) \tag{37}
\end{align*}
$$

There is still a length $x_{0}$ in the rolling direction when the workpiece exceeds the lowest point of the upper roll in the deformation zone. The mean value of the contact angle $\alpha_{m}$ in the downward rolling is similarly obtained using the integral mean value theorem

$$
\begin{equation*}
\alpha_{m}=\frac{\int_{-x_{0}}^{l} \alpha \mathrm{~d} x}{l+x_{0}}=\frac{\theta \sin \theta+\cos \theta-\alpha_{0} \sin \alpha_{0}-\cos \alpha_{0}}{\sin \theta+\sin \alpha_{0}} . \tag{38}
\end{equation*}
$$

The friction power $\dot{W}_{\mathrm{df}}$ is

$$
\begin{align*}
& \dot{W}_{\mathrm{df}}=4 m k w\left(\int_{-x_{0}}^{l}\left|\Delta v_{x}\right| \cos \alpha \sec \alpha \mathrm{d} x+\int_{-x_{0}}^{l}\left|\Delta v_{y}\right| \sin \alpha \sec \alpha \mathrm{d} x\right) \\
& =4 m k w\left\langle v_{R} R\left(\theta-2 \alpha_{n}-\frac{\sin 2 \alpha_{0}}{2}\right)+V_{y} R\left(\ln \frac{\cos ^{2} \alpha_{n} \cos \alpha_{0}}{\cos \theta}\right)\right. \\
& -V_{y} R\left(2 \cos \alpha_{n}-\cos \theta-\cos \alpha_{0}\right)+\frac{2\left(U / w-V_{y} R \theta+V_{y} R \alpha_{m}\right)}{M \sqrt{M^{2}-1}} \\
& \left\{\left(2 M^{2}-1\right) \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{0}}{2}\right)\right]-\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\theta}{2}\right)\right]\right.  \tag{39}\\
& -\alpha_{0} M \sqrt{M^{2}-1}+2 \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{n}}{2}\right)\right] \\
& \left.\left.+\frac{\sqrt{M^{2}-1}}{2} \ln \frac{\tan ^{2}\left(\pi / 4+\alpha_{n} / 2\right) \tan \left(\pi / 4-\alpha_{0} / 2\right)}{\tan (\pi / 4+\theta / 2)}\right\}\right\rangle .
\end{align*}
$$

using the MY liner yield criterion. Hence, the mean value of $\cos \alpha$ in deformation zone is obtained as

$$
\begin{equation*}
\overline{\cos \alpha}=\frac{\int_{-\alpha_{0}}^{\theta} \cos \alpha \mathrm{d} \alpha}{\theta+\alpha_{0}}=\frac{\sin \theta+\sin \alpha_{0}}{\theta+\alpha_{0}} \tag{35}
\end{equation*}
$$

The tension power $\dot{W}_{\mathrm{dT}}$ is

$$
\begin{equation*}
\dot{W}_{\mathrm{dT}}=4 U\left(\sigma_{\mathrm{b}}-\sigma_{\mathrm{f}}\right) \tag{40}
\end{equation*}
$$

According to Eqs. (36), (37), (39) and (40), the total power functional model in the downward rolling is

The internal plastic deformation power $\dot{W}_{\text {di }}$ becomes

$$
\begin{align*}
& \dot{W}_{\mathrm{di}}=\frac{32}{7} \sigma_{\mathrm{s}} \int_{-x_{0}}^{l} \int_{0}^{w} \int_{0}^{h_{x}}\left(\frac{-h_{0} v_{0}+V_{y} R \theta-V_{y} R \alpha}{h_{x}^{2}} h_{x}^{\prime}+\frac{V_{y} R}{x h_{x}} h_{x}^{\prime}\right) \mathrm{d} x \mathrm{~d} y \mathrm{~d} z \\
& =\frac{32}{7} \sigma_{\mathrm{s}} U \ln \frac{h_{0}}{h_{1}}+\frac{32}{7} \sigma_{\mathrm{s}} w V_{y} R\left\{\theta \ln \frac{h_{1}(M-\cos \theta)}{h_{0}}\right.  \tag{36}\\
& \left.+\alpha_{0} \ln \left(M-\cos \alpha_{0}\right)-\left(\theta+\alpha_{0}\right)\left[1+\ln \left(M-\frac{\sin \theta+\sin \alpha_{0}}{\theta+\alpha_{0}}\right)\right]\right\}
\end{align*}
$$



Figure 4 Flowchart of the calculation in variable gauge rolling

$$
\begin{equation*}
J_{\mathrm{d}}^{*}=\dot{W}_{\mathrm{di}}+\dot{W}_{\mathrm{ds}}+\dot{W}_{\mathrm{df}}+\dot{W}_{\mathrm{dT}} \tag{41}
\end{equation*}
$$

By solving the differential of the total power functional in Eq. (41) with respect to the arbitrary variable $\alpha_{n}$ and making it equal to zero, the following equation can be obtained as

$$
\begin{equation*}
\frac{\mathrm{d} J_{\mathrm{d}}^{*}}{\mathrm{~d} \alpha_{n}}=\frac{\mathrm{d} \dot{W}_{\mathrm{di}}}{\mathrm{~d} \alpha_{n}}+\frac{\mathrm{d} \dot{W}_{\mathrm{ds}}}{\mathrm{~d} \alpha_{n}}+\frac{\mathrm{d} \dot{W}_{\mathrm{df}}}{\mathrm{~d} \alpha_{n}}+\frac{\mathrm{d} \dot{W}_{\mathrm{dT}}}{\mathrm{~d} \alpha_{n}}=0 \tag{42}
\end{equation*}
$$

where

$$
\begin{equation*}
\frac{\mathrm{d} \dot{W}_{\mathrm{di}}}{\mathrm{~d} \alpha_{n}}=\frac{32}{7} \sigma_{\mathrm{s}} N \ln \frac{h_{0}}{h_{1}} \tag{43}
\end{equation*}
$$



Figure 5 Variation of roll gap with time in MAS rolling


Figure 6 Comparison of roll separating force for theoretical and measured values in MAS rolling

$$
\begin{aligned}
& \text { Rolling direction } \\
& \begin{array}{l}
\text { Unit: mm } \\
\text { Clamping Thick Downward } \\
\text { area } \quad \text { Thin } \\
\text { area rolling area } \\
\text { area }
\end{array} \text { Upward Thick Clamping } \\
& \text { rolling area area area }
\end{aligned}
$$

Figure 7 Geometry dimension of the workpiece after rolling

$$
\begin{equation*}
\frac{\mathrm{d} \dot{W}_{\mathrm{ds}}}{\mathrm{~d} \alpha_{n}}=2 k N\left(\tan \theta+\tan \alpha_{0}\right) \tag{44}
\end{equation*}
$$

and the roll separating force in the elastic zone is calculated based on our previous cold rolling research [30]. The calculation flow chart is depicted in Figure 4.

$$
\begin{align*}
& \frac{\mathrm{d} \dot{W}_{\mathrm{df}}}{\mathrm{~d} \alpha_{n}}=8 m k w\left\langle-V_{y} R \tan \alpha_{n}+V_{y} R \sin \alpha_{n}+\frac{U / w-V_{y} R \theta+V_{y} R \alpha_{m}}{\cos \alpha_{n}\left(M-\cos \alpha_{n}\right)}\right. \\
& -v_{R} R+\frac{N}{w M \sqrt{M^{2}-1}}\left\{\left(2 M^{2}-1\right) \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{0}}{2}\right)\right]-\alpha_{0} M .\right.  \tag{45}\\
& \sqrt{M^{2}-1}-\arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\theta}{2}\right)\right]+2 \arctan \left[\sqrt{\frac{M+1}{M-1}} \tan \left(\frac{\alpha_{n}}{2}\right)\right] \\
& \left.\left.+\frac{\sqrt{M^{2}-1}}{2} \ln \frac{\tan ^{2}\left(\pi / 4+\alpha_{n} / 2\right) \tan \left(\pi / 4-\alpha_{0} / 2\right)}{\tan (\pi / 4+\theta / 2)}\right\}\right\rangle,
\end{align*}
$$

$$
\begin{equation*}
\frac{\mathrm{d} \dot{W}_{\mathrm{dT}}}{\mathrm{~d} \alpha_{n}}=4 N\left(\sigma_{\mathrm{b}}-\sigma_{\mathrm{f}}\right) \tag{46}
\end{equation*}
$$

## 5 Results and Discussions

The minimum angle $\alpha_{n}$ (neutral angle) and the minimum value of the total power $J_{\text {min }}^{*}$ are obtained in different conditions of the upward and downward rolling based on Eqs. (27), (28), (41) and (42). The analytical models of the roll torque $M$, roll separating force $F$, and stress state effect coefficient $n_{\sigma}$ could be calculated as Eq.(47) [29].

$$
\begin{equation*}
M=\frac{R_{0} J_{\min }^{*}}{2 v_{R}}, F=\frac{M}{\chi l_{\text {total }}}, n_{\sigma}=\frac{F}{4 w l_{\text {total }} k} . \tag{47}
\end{equation*}
$$

The work roll radius is obviously flattened for the cold variable gauge rolling with front and back tensions, such as TRB rolling. Hitchcock's roll elastic flattening model with consideration of front and back tensions is adopted,


Figure 9 Variations of the bite angle, neutral angle, and location of neutral points in upward rolling


Figure 8 Comparison of roll separating force in a TRB


Figure $\mathbf{1 0}$ Variations of the bite angle, neutral angle, and location neutral points in downward rolling


Figure 11 Variation of powers in upward rolling

The mathematical model used in this research (present model) is certified using the measured data of the MAS rolling process in a medium plate hot rolling plant. The rolling technology is depicted in Figure 1(a). The workpiece material is Q345, the thickness is 0.182 m , and the width is 2.611 m . The roll diameter is 0.946 m , and the roll circumferential speed is $0.951 \mathrm{~m} / \mathrm{s}$. The change of the roll gap with time is depicted in Figure 5. The downward rolling process is performed at the head of the workpiece. Then the traditional equal thickness rolling process is carried out at high speed. Finally, the upward rolling process is executed near the end of the workpiece. The roll separating force in this process that is calculated with


Figure 13 Influence of friction factor on stress state coefficient in upward rolling
the present model is contrasted with measured values, as depicted in Figure 6. The maximum deviation is less than $5 \%$, except for the fact that the predicted roll separating force is larger at the starting point of the downward rolling and the endpoint of the upward rolling.

The predicted roll separating force accuracy of the present model in the TRB cold rolling process is verified using Zhang's research [21], in which a high strength micro-alloyed steel CR340 workpiece was rolled by a 450 mm four-high cold rolling mill with hydraulic tension in the laboratory. The workpiece is rolled as shown in Figure 7 (the unit is mm ). The original thickness for the


Figure 14 Influence of friction factor on stress state coefficient in downward rolling
workpiece is 2 mm and the front and back tensions are 40 kN .
The roll separating force computed with the present model is contrasted with the measured data in a rolling experiment and the values calculated using the slab method based on the idea of solving the Karman differential equation, as shown in Figure 8. The rolls are flattened and bounced severely due to the large material deformation resistance of the plate in cold rolling. The actual roll gap has to be less than the target roll gap to ensure the finished product size, as indicated in Figure 7. Therefore, the roll separating force is not zero at the beginning of the downward rolling and the end of the upward rolling. The roll separating force computed by the present model is larger than the computed results using the slab method, and it is slightly larger than the measured data. However, the compared results show good agreement, and the present model meets the requirements and accuracy of the roll separating force setting value. Therefore, the mathematical model used in this research can be exploited to forecast the roll separating force and to research variable gauge rolling.
Figures 9 and 10 show the variations of the bite angle, neutral angles, and locations of neutral points with time for upward and downward rolling, using the MAS rolling process parameters illustrated in Figure 5. The variation of the neutral angle in variable gauge rolling is compared with that of the slab method. The neutral angle calculated by the present model is slightly less than that calculated with the slab method because the calculation formula of the neutral angle in the slab method is based on the assumption that the unit pressure is uniformly distributed along the contact arc. The bite and neutral angles decrease, and the decrease rate of the bite angle is more remarkable than that of the neutral angle in upward rolling due to the reduction rate of workpiece gradually decreasing. On the contrary, the bite and neutral angles increase, and the increase rate of bite angle is more significant than that of neutral angle in downward rolling due to the gradual decrease of the reduction rate of the workpiece. According to Figures 9 and 10, the neutral point's location moves toward the inlet of the zone in upward rolling and moves toward the exit of the zone in downward rolling.
The variations of the internal plastic deformation power $\dot{W}_{\mathrm{i}}$, shear power $\dot{W}_{\mathrm{s}}$, and friction power $\dot{W}_{\mathrm{f}}$ with time in upward and downward rolling are depicted in Figures 11 and 12. It can be seen from the figures that the internal plastic deformation and shear powers make up a larger proportion of the total power. In comparison, the friction power makes up a minor proportion. They all increase with time in the upward rolling but decrease with time in the downward rolling. Moreover, the internal plastic
deformation power obviously varies with time, while the friction power only changes slightly.
Figures 13 and 14 show that the stress state effect coefficient $n_{\sigma}$ increases linearly with the increase of the rolling time in upward rolling, while it decreases with the increase of the rolling time in downward rolling using the MAS rolling process parameters displayed in Figure 5. This is because the shape factor of the rolled piece is less than 1 ; i.e., the ratio of the projected length to the average thickness is less than 1 in the deformation zone. Furthermore, for the deformation of the workpiece during variable gauge rolling, the effect of the external zone should be considered. Deformation not only occurs in the contact zone should be considered but also in the external zone. It is evident from Figures 11 and 12 that the shear power is greater than the friction power; that is to say, the influence of the external zone is the primary influence. Therefore, when the shape factor of the workpiece is small, the shear power is considerable. The external force must be increased to deform the workpiece, and the stress state effect coefficient must correspondingly increase. Figures 13 and 14 also show that the stress state effect coefficient $n_{\sigma}$ obviously increases linearly with the increase of the friction factor $m$ in the upward and downward rolling.

## 6 Conclusions

(1) The velocity and the strain rate fields in variable gauge rolling are built with consideration of the influence of the upward or downward motion of the rolls, with this field satisfying kinematically admissible conditions. The analytical models of the roll separating force, roll torque, and stress state effect coefficient are obtained.
(2) Comparisons among the predicted roll separating forces and measured values in MAS rolling in a medium plate hot rolling plant and a TRB rolling experiment show that the analytical model exhibits excellent prediction accuracy. The present model can, therefore, be utilized to research variable gauge rolling.
(3) The bite and neutral angles decrease in the upward rolling and increase in the downward rolling. The neutral point goes toward the inlet of the zone in upward rolling and moves toward the exit of the zone in downward rolling. They are the essential parameters of the micro-tracing strategy, which determines the length of the transition zone of the variable thickness plates.
(4) The stress state effect coefficient increases in the upward rolling and decreases in the downward rolling. The friction factor influences the stress state
effect coefficient noticeably. The thickness control strategy is developed based on the force parameters and the rolling mill bounce equation, which determines the thickness of the transition zone of the variable thickness plates.

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## Author Contributions

YL, QH and DZ were in charge of the whole trial; ZW and TW wrote the manuscript; JS and XZ assisted with result analyses. All authors read and approved the final manuscript.

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## Competing Interests

The authors declare no competing financial interests.

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